It is *Not* **Possible to Design a Rotationally Invariant Sampling Scheme for DT-MRI**

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INTRODUCTION: The choice of gradient sampling orientations for DT-MRI experiments has exercised many groups in the field. Of particular interest is the design of sampling schemes that will ensure that the error / variance in tensor estimates is independent of the relative orientation of the tensor to the reference frame established by the sampling vectors, which others have termed 'statistical rotational invariance' or '*SRI*^{'1,2}. Previous work has identified a relationship between the condition number of the quadratic encoding matrix, formed from the gradient sampling orientations, and the variance in the estimated diffusion tensor³. Batchelor *et al.*4 showed that certain sampling schemes (e.g., the dual-gradient scheme) had rotationally *variant* condition numbers while, those derived from the vertices of an icosahedron had rotationally *invariant* condition numbers. Thus, schemes based on vectors pointing to the vertices of an icosahedron (or tessellated icosahedrons) cause initial excitement. However, Jones² showed that, in discord with this theory, not all icosahedral schemes are the same – and that rotationally invariant condition number is a necessary but insufficient requirement for rotational invariance. A more general framework has been proposed^{1,5} – which provides a template for sampling schemes in terms of the fourth order precision matrix. Previous theoretical considerations suggested that, under the *linear* framework, a statistically rotationally invariant sampling scheme could not be designed, except for the trivial case of isotropic tensors⁵. Here, we address an important and long outstanding question in the DT-MRI literature, i.e. In the limit of an infinite number of sampling orientations, can one design a sampling scheme which is statistically rotationally invariant?

THEORY: By combining results from previous works^{1,5,6}, we can show how the precision matrix, M', in estimation of the diffusion tensor, **D**, can be expressed in terms of the elements of the m unit sampling vectors, **g**, the b-values, b_m and resultant **B** matrices, \mathbf{B}_m (where $\mathbf{B}_m = b_m \mathbf{g}_m^{-1} \mathbf{g}_m$), via Eq. [1], where $\alpha = 0$ and 1 for nonlinear and linear regression, respectively:

$$
\mathbf{M} = \begin{bmatrix} \frac{b_m^2 g_{\infty}^4}{\alpha + \exp(2Tr(\mathbf{B}_m \mathbf{D})} & \frac{b_m^2 g_{\infty}^2 g_{\infty}^2 g_{\infty}^2}{\alpha + \exp(2Tr(\mathbf{B}_m \mathbf{D}))} & 2 \frac{b_m^2 g_{\infty}^2 g_{\infty}^2 g_{\infty}^2 g_{\infty}^2}{\alpha + \exp(2Tr(\mathbf{B}_m \mathbf{D}))} & 2 \frac{b_m^2 g_{\infty}^2 g_{\infty}^2 g_{\infty}^2}{\alpha + \exp(2Tr(\mathbf{B}_m \mathbf{D}))} & 2 \frac{b_m^2 g_{\infty}^2 g_{\infty}^2 g_{\infty}^2}{\alpha + \exp(2Tr(\mathbf{B}_m \mathbf{D}))} & \frac{b_m^2 g_{\infty}^2 g_{\infty}^2 g_{\infty}^2 g_{\infty}^2}{\alpha + \exp(2Tr(\mathbf{B}_m \mathbf{D}))} & 2 \frac{b_m^2 g_{\infty}^2 g_{\infty}^2 g_{\infty}^2}{\alpha + \exp(2Tr(\mathbf{B}_m \mathbf{D}))} & 2 \frac{b_m^2 g_{\infty}^2 g_{\infty}^2 g_{\infty}^2}{\alpha + \exp(2Tr(\mathbf{B}_m \mathbf{D}))} & 2 \frac{b_m^2 g_{\infty}^2 g_{\infty}^2 g_{\infty}^2}{\alpha + \exp(2Tr(\mathbf{B}_m \mathbf{D}))} & 2 \frac{b_m^2 g_{\infty}^2 g_{\infty}^2 g_{\infty}^2 g_{\infty}^2}{\alpha + \exp(2Tr(\mathbf{B}_m \mathbf{D}))} & 2 \frac{b_m^2 g_{\infty}^2 g_{\infty}^2 g_{\infty}^2 g_{\infty}^2 g_{\infty}^2 g_{\infty}^2}{\alpha + \exp(2Tr(\mathbf{B}_m \mathbf{D}))} & 2 \frac{b_m^2 g_{\infty}^2 g_{\infty}^2 g_{\infty}^2 g_{\infty}^2 g_{\infty}^2 g_{\infty}^2}{\alpha + \exp(2Tr(\mathbf{B}_m \mathbf{D}))} &
$$

It has been shown that for the precision matrix to be statistically rotationally invariant, it should take the following general form^{1,7}:

$$
\mathbf{In Eq. [1], } \mathbf{M'}_{4,4} = 4\mathbf{M'}_{1,2}, \mathbf{M'}_{5,5} = 4\mathbf{M'}_{1,3} \text{ and } \mathbf{M'}_{6,6} = 4\mathbf{M'}_{2,3}, \text{ which on comparison with Eq. [2] means that } \lambda \text{ must be equal to } \mu_1. \text{ To conform to the prescription in Eq.}
$$
\n
$$
\begin{bmatrix}\n\lambda + 2\mu & \lambda & \lambda & 0 & 0 & 0 \\
\lambda & \lambda + 2\mu & \lambda & 0 & 0 & 0 \\
0 & 0 & 0 & 4\mu & 0 & 0 \\
0 & 0 & 0 & 0 & 4\mu\n\end{bmatrix}
$$
\n
$$
\begin{bmatrix}\n2 \\
\lambda & \lambda + 2\mu & \lambda & 0 & 0 \\
0 & 0 & 0 & 4\mu & 0 \\
0 & 0 & 0 & 0 & 4\mu\n\end{bmatrix}
$$
\n
$$
\begin{bmatrix}\n2 \\
\lambda\n\end{bmatrix}
$$

[2], *all* elements of M' should conform for *all* possible tensors. Therefore, if we can find just one example of non-conformance, we have shown that the design is not *SRI*. To make progress, we examine just the non-zero elements of M', which contain only even powers of the elements of sampling vector **g**. Thus, we define an even function of these elements, as: κ (g $_p$) = (g $_{x_p}$)²¹(g $_{x_p}$)^{2m} (g $_{z_p}$)^{2m}. Summing over an *infinite* number of directions, is the same as integrating over all angles – so we re-

parameterize **g** in spherical co-ordinates, $g = \left[\sin\theta\cos\phi, \sin\theta\sin\phi, \cos\phi\right]$, thus the form for the infinite sums in Eq. [1] becomes

$$
\langle f \rangle = \frac{b^2}{4\pi} \int_{\phi=0}^{\pi} \int_{\theta=0}^{\pi} \exp(-2Tr(\mathbf{BD})) \kappa(\mathbf{g}) \sin(\theta) d\theta d\phi
$$
 [3]

If we assume **D** is aligned with principal lab frame, (i.e., $\mathbf{D}_{ii} = \delta_{ii} \lambda_i$, where δ_{ii} is the Dirac delta function, and λ_i are eigenvalues), then

$$
\langle f \rangle = \frac{b^2}{4\pi} \int_{\phi=0}^{\pi} (g_X)^{2l} (g_Y)^{2m} (g_Z)^{2n} \exp(-(s_1 g_X^2 + s_2 g_Y^2 + s_3 g_Z^2)) \sin(\theta) d\theta d\phi
$$
 [4]

where $s_i = 2b\lambda_i$, for $i=1,2,3$.

First, we consider the trivial case of isotropic tensors, i.e., $\lambda_i = \lambda$. The exponent in Eq.[4] can be factored out as a scalar multiplier, and the result is the sum of an infinite number of sampling orientations which Batchelor et al.⁴ have previously shown takes the form in Eq. [2]. For a more general case, however, the evaluation of Eq.[4] becomes more complex. Space prohibits detailed working, (utilizing Mathematica), but the end result is:

$$
\left\langle f(l,m,n,s_1,s_2,s_3) \right\rangle = \frac{b^2}{2\sqrt{\pi}} e^{-s_3} \frac{\Gamma(1+l+m)\Gamma(\frac{1}{2}+n)}{\Gamma(\frac{3}{2}+l+m+n)} \sum_{k=0}^{\infty} \left(\frac{(1+l+m)_k}{k!\left(\frac{3}{2}+l+m+n\right)_k} \sum_{q=0}^k \left(\binom{k}{q} (s_3-s_1)^{k-q} (s_1-s_2)^q \frac{\Gamma(\frac{1}{2}+m+q)}{\Gamma(1+m+q)} {}_2F_1(-2l;\frac{1}{2}+m+q;1+2m+2q;2) \right) \right) \tag{5}
$$

where $\sum_{i=0}^{N} \frac{(-1)^{k+1} \sum_{j=0}^{N} (k-1)^{j}}{k! (k-1)(k-2)!}$
where $\sum_{i=0}^{N} \frac{(-1)^{k+1} (k-1)^{j}}{k! (k-1)(k-1)}$ and $\sum_{i=0}^{N} \frac{(-1)^{i-1} (k-1)^{j}}{k! (k-1)(k-1)}$ and $\sum_{i=0}^{N} \frac{(-1)^{i-1} (k-1)^{j}}{k! (k-1)(k-1)}$ is the Kummer

 $(a)_n = a(a+1)(a+2)\cdots(a+(n-1))$ is the Pochhammer function. We now assume a general tensor such that $(s_1, s_2, s_3) = (2b\lambda_1, 2b\lambda_2, 2b\lambda_3) = (0.1, 0.01, 0.001)$ and evaluate the leading three diagonal elements of M', $(M_{1,1} M_{2,2}$ and $M_{3,3}$) using Eq. [5]. The results are $f(2,0,0, 0.1, 0.01, 0.001) = 0.195056...$; $f(0,2,0,0.1,0.01,0.001) = 0.195754...$ and $f(0,0,2,0.1,0.01,0.001) = 0.190843...$ This indicates that the leading terms are not equal and therefore the prescription given in Eq. [2] is violated.

CONCLUSION: In the limit of an infinite number of sampling vectors, the precision matrix is rotationally invariant *only* for the trivial case of isotropic tensors – a result previously outlined elsewhere⁵. Even in this infinite limit, however, for a given anisotropic tensor, the leading 3 diagonal terms of M' are unequal and hence the condition laid out in Eq. [2] is violated. As discussed earlier, just one example of non-conformance to Eq. [2] is sufficient to conclude that, despite claims by some groups to the contrary, it is indeed NOT possible to design a statistically rotationally invariant sampling scheme.

REFERENCES: 1. Basser et al. *TMI* 2003 ; 22 :785- ; **2.** Jones *MRM* 2004; 51:807-; **3.** Skare et al. *J Magn Reson* 2000; 147: 340- ; **4.** Batchelor et al. *MRM* 2003; 49:1143-; **5.** Jones *Proc ISMRM* 2003, p. 2118. **6.** Koay et al. *J Magn Reson* 2006; 182:115-; **7.** Hext *Biometrika* 1963; 50:353-;